

## 6.632 Solution to Problem Set 9

### Solution P9.1

$$\bar{H}(\bar{r}) = (\hat{\theta} \cos \phi - \hat{\phi} \cos \theta \sin \phi) \frac{ik}{\pi r k_x k_y} e^{ikr} \sin \frac{k_x D_x}{2} \sin \frac{k_y D_y}{2}.$$

### Solution P9.2

The first zero of  $\sin(kxw/2z)$  occurs at  $kxw/2z = \pi$  which yields  $2\pi w\theta/2\lambda = \pi$ . Consequently  $\theta = \lambda/w$ .

For the  $N$ th side lobe,  $kxw/2z = kw\theta/2 = N\pi$ , which yields  $\theta_N = N\lambda/w$ . Based on the assumption  $\theta \approx x/z < 1$ , we thus deduce that  $N < w/\lambda$ .

### Solution P9.3

(a) For a biaxial medium,

$$\bar{\bar{\epsilon}} = \begin{bmatrix} \epsilon_x & 0 & 0 \\ 0 & \epsilon_y & 0 \\ 0 & 0 & \epsilon_z \end{bmatrix}$$

Since  $\bar{\bar{\epsilon}}^t = \bar{\bar{\epsilon}}$ , therefore the medium is reciprocal.

(b) For a moving biaxial medium,

$$\bar{\bar{C}}_{EB} = \gamma^2 \begin{bmatrix} c\epsilon'_x - \frac{\beta^2}{c\mu'} & 0 & 0 & 0 & \beta \left( -c\epsilon'_x + \frac{1}{c\mu'} \right) & 0 \\ 0 & c\epsilon'_y - \frac{\beta^2}{c\mu'} & 0 & \beta \left( c\epsilon'_y - \frac{1}{c\mu'} \right) & 0 & 0 \\ 0 & 0 & \frac{c\epsilon'_z}{\gamma^2} & 0 & 0 & 0 \\ 0 & \beta \left( -c\epsilon'_y + \frac{1}{c\mu'} \right) & 0 & -c\epsilon'_y \beta^2 + \frac{1}{c\mu'} & 0 & 0 \\ \beta \left( c\epsilon'_x - \frac{1}{c\mu'} \right) & 0 & 0 & 0 & -\beta^2 c\epsilon'_x + \frac{1}{c\mu'} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{\gamma^2 c\mu'} \end{bmatrix} \quad (1)$$

Since

$$\begin{bmatrix} c\bar{D} \\ \bar{H} \end{bmatrix} = \begin{bmatrix} \bar{P} \\ \bar{M} \end{bmatrix} \begin{bmatrix} \bar{L} \\ \bar{Q} \end{bmatrix} \begin{bmatrix} \bar{E} \\ c\bar{B} \end{bmatrix}$$

$$\bar{\bar{C}}_{EH} = \begin{bmatrix} \bar{\bar{\epsilon}} & \bar{\bar{\xi}} \\ \bar{\bar{\zeta}} & \bar{\bar{\mu}} \end{bmatrix} = \frac{1}{c} \begin{bmatrix} \bar{P} - \bar{L} \cdot \bar{Q}^{-1} \cdot \bar{M} & \bar{L} \cdot \bar{Q}^{-1} \\ -\bar{Q}^{-1} \cdot \bar{M} & \bar{Q}^{-1} \end{bmatrix} \quad (2)$$

We observe that  $\bar{\bar{\xi}}$  and  $\bar{\bar{\zeta}}$  are not imaginary since there are no imaginary elements in (1). Thus the medium is nonreciprocal. The complementary medium is characterized by

$$\begin{aligned} \bar{\bar{\mu}}^c &= \bar{\bar{\mu}}^t; & \bar{\bar{\epsilon}}^c &= \bar{\bar{\epsilon}}^t; \\ \bar{\bar{\zeta}}^c &= -\bar{\bar{\zeta}}^t; & \bar{\bar{\xi}}^c &= -\bar{\bar{\xi}}^t. \end{aligned}$$

(c) A biisotropic medium with imaginary  $\chi$  is reciprocal.

(d) For a biisotropic medium with a real  $\chi$ .

$$\begin{aligned}\overline{D} &= \epsilon \overline{E} + \xi \overline{H} \\ \overline{B} &= -\xi \overline{E} + \mu \overline{H}\end{aligned}$$

the medium is reciprocal but lossy.

(e) For a ferrite in a dc magnetic field:

$$\overline{\overline{\mu}} = \begin{bmatrix} \mu_1 & -i\mu_2 & 0 \\ i\mu_2 & \mu_1 & 0 \\ 0 & 0 & \mu_0 \end{bmatrix}$$

and the medium is not reciprocal. The complementary medium is

$$\overline{\overline{\mu}}^c = \overline{\overline{\mu}}^t = \begin{bmatrix} \mu_1 & i\mu_2 & 0 \\ -i\mu_2 & \mu_1 & 0 \\ 0 & 0 & \mu_0 \end{bmatrix}$$